

1. [G 2.6] Find the electric field a distance z above the center of a flat circular disk of radius R (see figure), which carries a uniform surface charge σ . What does your formula give in the limit $R \rightarrow \infty$? Also check the case $z \gg R$.

Break it into rings of radius r , and thickness dr , and use Prob. 2.5 to express the field of each ring. Total charge of a ring is $\sigma \cdot 2\pi r \cdot dr = \lambda \cdot 2\pi r$, so $\lambda = \sigma dr$ is the "line charge" of each ring.

$$E_{\text{ring}} = \frac{1}{4\pi\epsilon_0} \frac{(\sigma dr)2\pi r z}{(r^2 + z^2)^{3/2}}; \quad E_{\text{disk}} = \frac{1}{4\pi\epsilon_0} 2\pi\sigma z \int_0^R \frac{r}{(r^2 + z^2)^{3/2}} dr.$$

$$\mathbf{E}_{\text{disk}} = \frac{1}{4\pi\epsilon_0} 2\pi\sigma z \left[\frac{1}{z} - \frac{1}{\sqrt{R^2 + z^2}} \right] \hat{\mathbf{z}}.$$

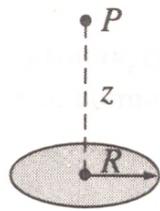
2. A spherical surface of radius R and center at origin carries a surface charge $\sigma(\theta, \phi) = \sigma_0 \cos\theta$. Find the electric field at z on z -axis. Treat the case $z < R$ (inside) as well as $z > R$ (outside). [Hint: Be sure to take the positive square root: $\sqrt{R^2 + z^2 - 2Rz} = (R - z)$ if $R > z$, but its $(z - R)$ if $R < z$.]

Let $\hat{\mathbf{r}} = z\hat{\mathbf{z}}$ and $\mathbf{r}' = R(\sin\theta' \cos\phi' \hat{\mathbf{x}} + \sin\theta' \sin\phi' \hat{\mathbf{y}} + \cos\theta' \hat{\mathbf{z}})$.

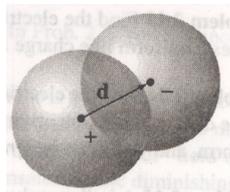
$$dS' = R^2 \sin\theta' d\theta' d\phi'$$

$$|\mathbf{r} - \mathbf{r}'| = (R^2 + z^2 - 2Rz \cos\theta')^{1/2}$$

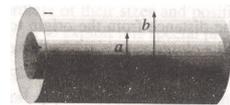
$$\begin{aligned} \mathbf{E}(z\hat{\mathbf{z}}) &= \frac{1}{4\pi\epsilon_0} \int_S \frac{\sigma(\mathbf{r}')(\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3} dS \\ &= \frac{\sigma_0}{4\pi\epsilon_0} \int_S \frac{\cos\theta' (-R \sin\theta' \cos\phi' \hat{\mathbf{x}} - R \sin\theta' \sin\phi' \hat{\mathbf{y}} + (z - R \cos\theta') \hat{\mathbf{z}})}{(R^2 + z^2 - 2Rz \cos\theta')^{3/2}} R^2 \sin\theta' d\theta' d\phi' \\ &= \frac{\sigma_0 R^2}{2\epsilon_0} \hat{\mathbf{z}} \int_{-1}^1 \frac{u(z - Ru)}{(R^2 + z^2 - 2Rzu)^{3/2}} du \\ &= \frac{2\sigma_0 R^3}{3\epsilon_0 z^3} \hat{\mathbf{z}} \quad z > R \end{aligned}$$



(a) Problem 1



(b) Problem 4



(c) Problem 5

Figure 1

$$= -\frac{\sigma_0}{3\epsilon_0}\hat{z} \quad z < R$$

Note: \mathbf{E} is constant inside sphere.

3. [G 2.9] Suppose the electric field in some region is found to be $\mathbf{E} = kr^3\hat{r}$, in spherical coordinates (k is some constant).

(a) Find the charge density ρ .

(b) Find the total charge contained in the sphere of radius R , centered at the origin. (Do it two different ways.)

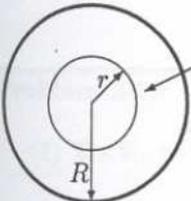
$$(a) \rho = \epsilon_0 \nabla \cdot \mathbf{E} = \epsilon_0 \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \cdot kr^3) = \epsilon_0 \frac{1}{r^2} k(5r^4) = \boxed{5\epsilon_0 kr^2}.$$

$$(b) \text{ By Gauss's law: } Q_{\text{enc}} = \epsilon_0 \oint \mathbf{E} \cdot d\mathbf{a} = \epsilon_0 (kR^3)(4\pi R^2) = \boxed{4\pi\epsilon_0 kR^5}.$$

$$\text{ By direct integration: } Q_{\text{enc}} = \int \rho d\tau = \int_0^R (5\epsilon_0 kr^2)(4\pi r^2 dr) = 20\pi\epsilon_0 k \int_0^R r^4 dr = 4\pi\epsilon_0 kR^5.$$

4. [G 2.12] Use Gauss's law to find the electric field inside a uniformly charged sphere (charge density ρ). [G 2.18] Two spheres each of radius R and carrying uniform charge densities $+\rho$ and $-\rho$, respectively, are placed so that they partially overlap (See Figure). Call the vector from the positive center to the negative center \mathbf{d} . Show that the field in the region of overlap is constant, and find its value.

[G 2.12]



Gaussian surface

$$\oint \mathbf{E} \cdot d\mathbf{a} = E \cdot 4\pi r^2 = \frac{1}{\epsilon_0} Q_{\text{enc}} = \frac{1}{\epsilon_0} \frac{4}{3}\pi r^3 \rho. \quad \text{So}$$

$$\boxed{\mathbf{E} = \frac{1}{3\epsilon_0} \rho r \hat{r}.}$$

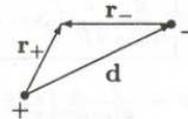
$$\text{ Since } Q_{\text{tot}} = \frac{4}{3}\pi R^3 \rho, \quad \mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{Q}{R^3} \hat{r}$$

[G 2.18]

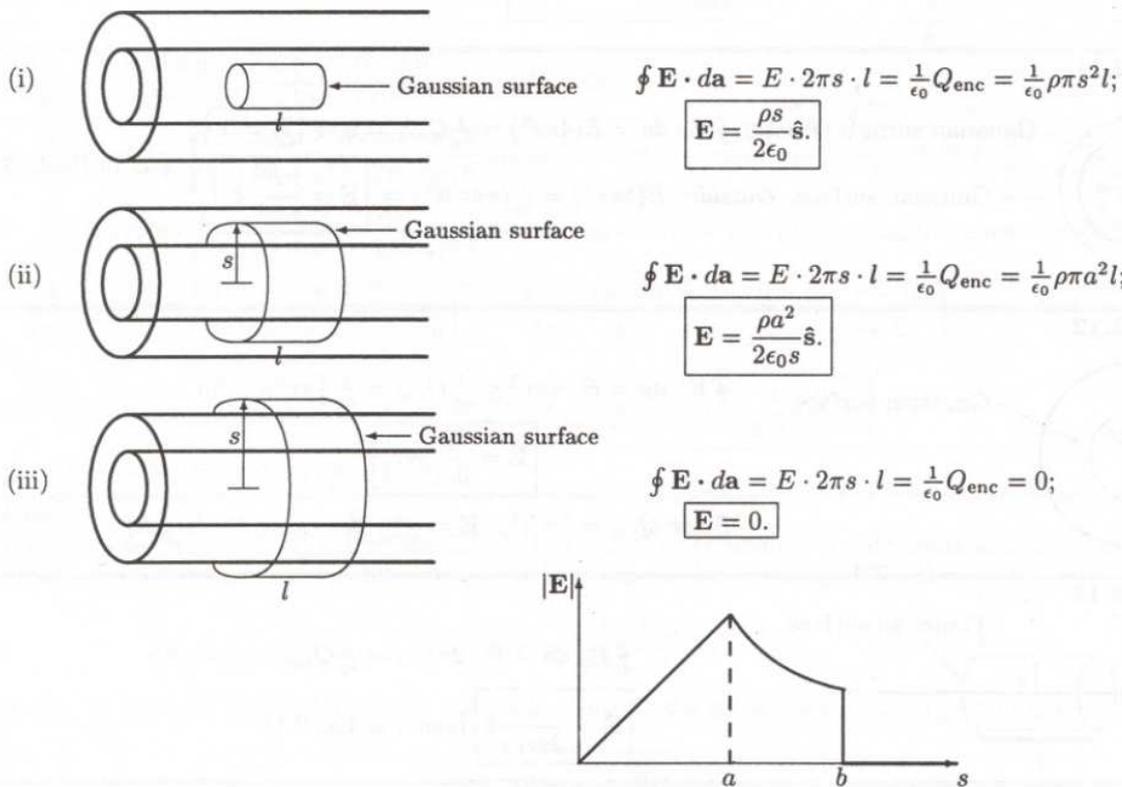
From Prob. 2.12, the field inside the positive sphere is $\mathbf{E}_+ = \frac{\rho}{3\epsilon_0} \mathbf{r}_+$, where \mathbf{r}_+ is the vector from the positive center to the point in question. Likewise, the field of the negative sphere is $-\frac{\rho}{3\epsilon_0} \mathbf{r}_-$. So the *total* field is

$$\mathbf{E} = \frac{\rho}{3\epsilon_0} (\mathbf{r}_+ - \mathbf{r}_-)$$

But (see diagram) $\mathbf{r}_+ - \mathbf{r}_- = \mathbf{d}$. So $\boxed{\mathbf{E} = \frac{\rho}{3\epsilon_0} \mathbf{d}.}$



5. [G 2.16] A long coaxial cable (see figure) carries a uniform volume charge density ρ on the inner cylinder (radius a), and a uniform surface charge density on the outer cylindrical shell (radius b). This surface charge is negative and of just the right magnitude so that the cable as a whole is electrically neutral. Find the electric field in each of the three regions: (a) inside the inner cylinder ($s < a$), (b) between the cylinders ($a < s < b$), (c) outside the cable ($s > b$). Plot $|\mathbf{E}|$ as a function of s .



6. [G 2.20] One of these is an impossible electrostatic field. Which one?

- (a) $\mathbf{E} = k [xy\hat{x} + 2yz\hat{y} + 3xz\hat{z}]$
 (b) $\mathbf{E} = k [y^2\hat{x} + (2xy + z^2)\hat{y} + 2yz\hat{z}]$.

Here k is a constant with the appropriate units. For the possible one, find the potential, using the origin as your reference point. Check your answers by computing ∇V . [Hint: You must select a specific path to integrate along. It does not matter what path you choose, since the answer is path-independent, but you simply cannot integrate unless you have a particular path in mind.]

$$(1) \nabla \times \mathbf{E}_1 = k \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ xy & 2yz & 3zx \end{vmatrix} = k [\hat{x}(0 - 2y) + \hat{y}(0 - 3z) + \hat{z}(0 - x)] \neq 0,$$

so \mathbf{E}_1 is an impossible electrostatic field.

$$(2) \nabla \times \mathbf{E}_2 = k \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ y^2 & 2xy + z^2 & 2yz \end{vmatrix} = k [\hat{x}(2z - 2z) + \hat{y}(0 - 0) + \hat{z}(2y - 2y)] = 0,$$

so \mathbf{E}_2 is a possible electrostatic field.

Let's go by the indicated path:

$$\mathbf{E} \cdot d\mathbf{l} = (y^2 dx + (2xy + z^2)dy + 2yz dz)k$$

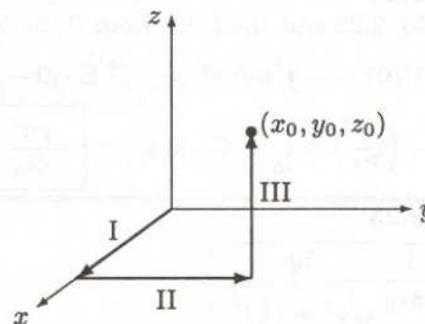
$$\text{Step I: } y = z = 0; dy = dz = 0. \mathbf{E} \cdot d\mathbf{l} = ky^2 dx = 0.$$

$$\text{Step II: } x = x_0, y : 0 \rightarrow y_0, z = 0. dx = dz = 0.$$

$$\mathbf{E} \cdot d\mathbf{l} = k(2xy + z^2)dy = 2kx_0y dy.$$

$$\int_{II} \mathbf{E} \cdot d\mathbf{l} = 2kx_0 \int_0^{y_0} y dy = kx_0y_0^2.$$

$$\text{Step III: } x = x_0, y = y_0, z : 0 \rightarrow z_0; dx = dy = 0.$$



$$\mathbf{E} \cdot d\mathbf{l} = 2kyz dz = 2ky_0z dz.$$

$$\int_{III} \mathbf{E} \cdot d\mathbf{l} = 2y_0k \int_0^{z_0} z dz = ky_0z_0^2.$$

$$V(x_0, y_0, z_0) = - \int_0^{(x_0, y_0, z_0)} \mathbf{E} \cdot d\mathbf{l} = -k(x_0y_0^2 + y_0z_0^2), \text{ or } \boxed{V(x, y, z) = -k(xy^2 + yz^2)}.$$

$$\text{Check: } -\nabla V = k \left[\frac{\partial}{\partial x}(xy^2 + yz^2) \hat{x} + \frac{\partial}{\partial y}(xy^2 + yz^2) \hat{y} + \frac{\partial}{\partial z}(xy^2 + yz^2) \hat{z} \right] = k[y^2 \hat{x} + (2xy + z^2) \hat{y} + 2yz \hat{z}] = \mathbf{E}.$$

7. [G 2.21] Find the potential inside and outside a uniformly charged solid sphere whose radius is R and whose total charge is q . Use infinity as your reference point. Compute the gradient of V in each region, and check that it yields the correct field. Sketch $V(r)$.

$$V(r) = -\int_{\infty}^r \mathbf{E} \cdot d\mathbf{l} \quad \begin{cases} \text{Outside the sphere } (r > R) : & \mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{\mathbf{r}}. \\ \text{Inside the sphere } (r < R) : & \mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} r \hat{\mathbf{r}}. \end{cases}$$

So for $r > R$: $V(r) = -\int_{\infty}^r \left(\frac{1}{4\pi\epsilon_0} \frac{q}{\bar{r}^2} \right) d\bar{r} = \frac{1}{4\pi\epsilon_0} q \left(\frac{1}{\bar{r}} \right) \Big|_{\infty}^r = \boxed{\frac{q}{4\pi\epsilon_0} \frac{1}{r}},$

and for $r < R$: $V(r) = -\int_{\infty}^R \left(\frac{1}{4\pi\epsilon_0} \frac{q}{\bar{r}^2} \right) d\bar{r} - \int_R^r \left(\frac{1}{4\pi\epsilon_0} \frac{q}{R^3} \bar{r} \right) d\bar{r} = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{R} - \frac{1}{R^3} \left(\frac{r^2 - R^2}{2} \right) \right]$
 $= \boxed{\frac{q}{4\pi\epsilon_0} \frac{1}{2R} \left(3 - \frac{r^2}{R^2} \right)}.$

When $r > R$, $\nabla V = \frac{q}{4\pi\epsilon_0} \frac{\partial}{\partial r} \left(\frac{1}{r} \right) \hat{\mathbf{r}} = -\frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$, so $\mathbf{E} = -\nabla V = \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}}$.

When $r < R$, $\nabla V = \frac{q}{4\pi\epsilon_0} \frac{1}{2R} \frac{\partial}{\partial r} \left(3 - \frac{r^2}{R^2} \right) \hat{\mathbf{r}} = \frac{q}{4\pi\epsilon_0} \frac{1}{2R} \left(-\frac{2r}{R^2} \right) \hat{\mathbf{r}} = -\frac{q}{4\pi\epsilon_0} \frac{r}{R^3} \hat{\mathbf{r}}$; so $\mathbf{E} = -\nabla V = \frac{1}{4\pi\epsilon_0} \frac{q}{R^3} r \hat{\mathbf{r}}$.

8. [G 2.22] Find the potential a distance s from an infinitely long straight wire that carries a uniform line charge λ . Compute the gradient of your potential, and check that it yields the correct field.

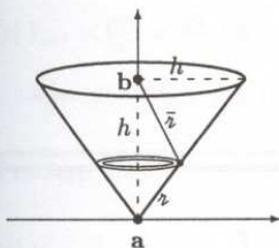
$\mathbf{E} = \frac{1}{4\pi\epsilon_0} \frac{2\lambda}{s} \hat{\mathbf{s}}$ (Prob. 2.13). In this case we cannot set the reference point at ∞ , since the charge itself extends to ∞ . Let's set it at $s = a$. Then

$$V(s) = -\int_a^s \left(\frac{1}{4\pi\epsilon_0} \frac{2\lambda}{\bar{s}} \right) d\bar{s} = \boxed{-\frac{1}{4\pi\epsilon_0} 2\lambda \ln \left(\frac{s}{a} \right)}.$$

(In this form it is clear why $a = \infty$ would be no good—likewise the other “natural” point, $a = 0$.)

$$\nabla V = -\frac{1}{4\pi\epsilon_0} 2\lambda \frac{\partial}{\partial s} \left(\ln \left(\frac{s}{a} \right) \right) \hat{\mathbf{s}} = -\frac{1}{4\pi\epsilon_0} 2\lambda \frac{1}{s} \hat{\mathbf{s}} = -\mathbf{E}.$$

9. [G 2.26] A conical surface (an empty ice-cream cone) carries a uniform surface charge σ . The height of the cone is h , as is the radius of the top. Find the potential difference between points **a** (the vertex) and **b** (the center of the top).



$$V(\mathbf{a}) = \frac{1}{4\pi\epsilon_0} \int_0^{\sqrt{2}h} \left(\frac{\sigma 2\pi r}{z} \right) dz = \frac{2\pi\sigma}{4\pi\epsilon_0} \frac{1}{\sqrt{2}} (\sqrt{2}h) = \frac{\sigma h}{2\epsilon_0}$$

(where $r = z/\sqrt{2}$)

$$V(\mathbf{b}) = \frac{1}{4\pi\epsilon_0} \int_0^{\sqrt{2}h} \left(\frac{\sigma 2\pi r}{\bar{z}} \right) dz, \quad \text{where } \bar{z} = \sqrt{h^2 + z^2} - \sqrt{2}hz.$$

$$= \frac{2\pi\sigma}{4\pi\epsilon_0} \frac{1}{\sqrt{2}} \int_0^{\sqrt{2}h} \frac{z}{\sqrt{h^2 + z^2} - \sqrt{2}hz} dz$$

$$= \frac{\sigma}{2\sqrt{2}\epsilon_0} \left[\sqrt{h^2 + z^2} - \sqrt{2}hz + \frac{h}{\sqrt{2}} \ln(2\sqrt{h^2 + z^2} - \sqrt{2}hz + 2z - \sqrt{2}h) \right]_0^{\sqrt{2}h}$$

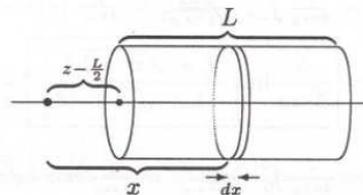
$$= \frac{\sigma}{2\sqrt{2}\epsilon_0} \left[h + \frac{h}{\sqrt{2}} \ln(2h + 2\sqrt{2}h - \sqrt{2}h) - h - \frac{h}{\sqrt{2}} \ln(2h - \sqrt{2}h) \right] = \frac{\sigma}{2\sqrt{2}\epsilon_0} \frac{h}{\sqrt{2}} \left[\ln(2h + \sqrt{2}h) - \ln(2h - \sqrt{2}h) \right]$$

$$= \frac{\sigma h}{4\epsilon_0} \ln \left(\frac{2 + \sqrt{2}}{2 - \sqrt{2}} \right) = \frac{\sigma h}{4\epsilon_0} \ln \left(\frac{(2 + \sqrt{2})^2}{2} \right) = \frac{\sigma h}{2\epsilon_0} \ln(1 + \sqrt{2}).$$

$$\therefore V(\mathbf{a}) - V(\mathbf{b}) = \frac{\sigma h}{2\epsilon_0} \left[1 - \ln(1 + \sqrt{2}) \right].$$

10. [G2.27] Find the potential on the axis of a uniformly charged solid cylinder, a distance z from the center. The length of the cylinder is L and its radius is R , and the charge density is ρ . Use your result to calculate the electric field at this point. (Assume that $z > L/2$.)

Cut the cylinder into slabs, as shown in the figure, and use result of Prob. 2.25c, with $z \rightarrow x$ and $\sigma \rightarrow \rho dx$:



$$V = \frac{\rho}{2\epsilon_0} \int_{z-L/2}^{z+L/2} (\sqrt{R^2 + x^2} - x) dx$$

$$= \frac{\rho}{2\epsilon_0} \frac{1}{2} [x\sqrt{R^2 + x^2} + R^2 \ln(x + \sqrt{R^2 + x^2}) - x^2]_{z-L/2}^{z+L/2}$$

$$= \frac{\rho}{4\epsilon_0} \left\{ (z + \frac{L}{2})\sqrt{R^2 + (z + \frac{L}{2})^2} - (z - \frac{L}{2})\sqrt{R^2 + (z - \frac{L}{2})^2} + R^2 \ln \left[\frac{z + \frac{L}{2} + \sqrt{R^2 + (z + \frac{L}{2})^2}}{z - \frac{L}{2} + \sqrt{R^2 + (z - \frac{L}{2})^2}} \right] - 2zL \right\}$$

(Note: $-(z + \frac{L}{2})^2 + (z - \frac{L}{2})^2 = -z^2 - zL - \frac{L^2}{4} + z^2 - zL + \frac{L^2}{4} = -2zL$.)

$$\begin{aligned} \mathbf{E} = -\nabla V = -\hat{\mathbf{z}} \frac{\partial V}{\partial z} = -\frac{\hat{\mathbf{z}}\rho}{4\epsilon_0} & \left\{ \sqrt{R^2 + \left(z + \frac{L}{2}\right)^2} + \frac{\left(z + \frac{L}{2}\right)^2}{\sqrt{R^2 + \left(z + \frac{L}{2}\right)^2}} - \sqrt{R^2 + \left(z - \frac{L}{2}\right)^2} - \frac{\left(z - \frac{L}{2}\right)^2}{\sqrt{R^2 + \left(z - \frac{L}{2}\right)^2}} \right. \\ & \left. + R^2 \left[\frac{1 + \frac{z + \frac{L}{2}}{\sqrt{R^2 + \left(z + \frac{L}{2}\right)^2}}}{z + \frac{L}{2} + \sqrt{R^2 + \left(z + \frac{L}{2}\right)^2}} - \frac{1 + \frac{z - \frac{L}{2}}{\sqrt{R^2 + \left(z - \frac{L}{2}\right)^2}}}{z - \frac{L}{2} + \sqrt{R^2 + \left(z - \frac{L}{2}\right)^2}} \right] - 2L \right\} \\ & \frac{1}{\sqrt{R^2 + \left(z + \frac{L}{2}\right)^2}} - \frac{1}{\sqrt{R^2 + \left(z - \frac{L}{2}\right)^2}} \end{aligned}$$

$$\mathbf{E} = -\frac{\hat{\mathbf{z}}\rho}{4\epsilon_0} \left\{ 2\sqrt{R^2 + \left(z + \frac{L}{2}\right)^2} - 2\sqrt{R^2 + \left(z - \frac{L}{2}\right)^2} - 2L \right\}$$

$$= \frac{\rho}{2\epsilon_0} \left[L - \sqrt{R^2 + \left(z + \frac{L}{2}\right)^2} + \sqrt{R^2 + \left(z - \frac{L}{2}\right)^2} \right] \hat{\mathbf{z}}.$$