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# Physics I

## *Lecture 4*

Charudatt Kadolkar

IIT Guwahati

# Kinetic Energy

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Kinetic Energy of a particle in motion is defined as

$$KE = \frac{1}{2}mv^2$$

where  $v$  is the instantaneous velocity of the particle.

In 2D, clearly

$$\begin{aligned}\frac{1}{2}mv^2 &= \frac{1}{2}m\mathbf{v} \cdot \mathbf{v} \\ &= \frac{1}{2}m(v_x^2 + v_y^2)\end{aligned}$$

Kinetic Energy changes as particle moves.

# Example

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A particle is moving on a circular path in horizontal plane.  $R = 1$  m,  $m = 1$  kg.  $\mu = 0.3$ . At  $t = 0$ ,  $\theta_0 = 0$ ,  $\omega_0 = 9$  rad/sec.

The retarding acceleration  $\alpha = \mu g = 3$  rad/sec<sup>2</sup>.

Velocity is given by  $R(\omega_0 - \alpha t) = (9 - 3t)$  m/s.

Kinetic Energy is then

$$\begin{aligned} KE(t) &= \frac{1}{2}m(v(t))^2 \\ &= \frac{1}{2}(9 - 3t)^2 \text{ J} \end{aligned}$$

Kinetic energy diminishes.

# Taylor Series

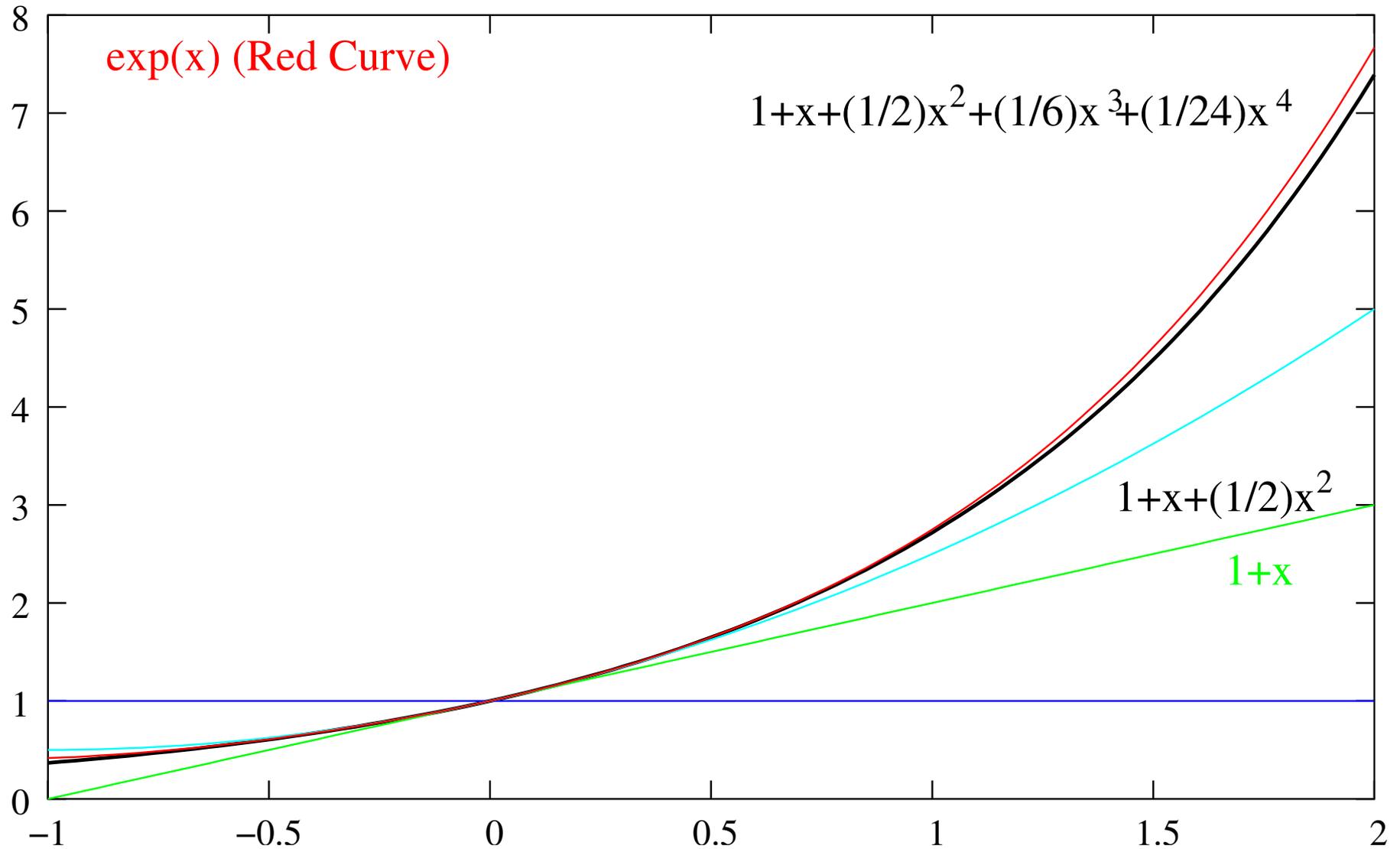
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Let  $f(x)$  be differentiable to any order. The Taylor theorem is

$$f(x + \delta x) = f(x) + f'(x)\delta x + \frac{1}{2!}f''(x)(\delta x)^2 + \dots + R_n(x)$$

- If the series is convergent, the remainder  $R_n(x)$  vanishes as  $n \rightarrow \infty$
- The function  $f$  can be approximated in the vicinity of  $x$  by polynomial in  $\delta x$  if  $\delta x$  is sufficiently small.

# Taylor Series

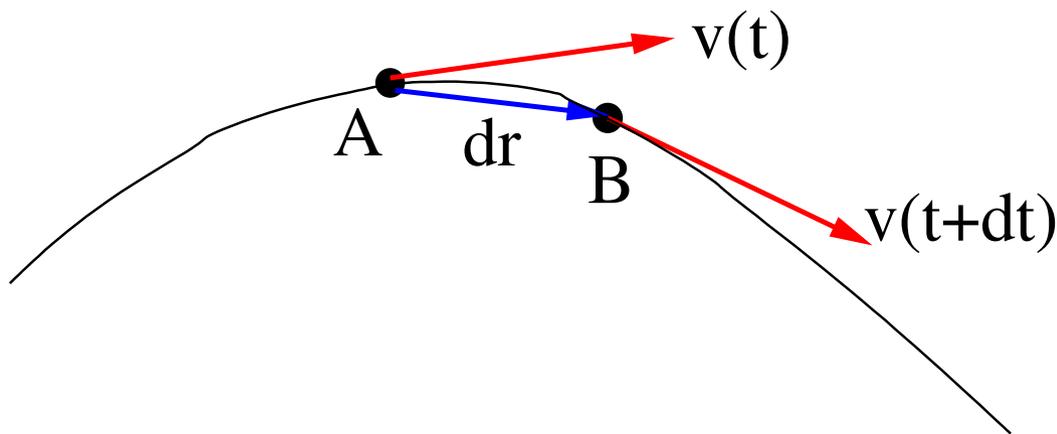


# Work

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A particle moves a little distance  $\Delta \mathbf{r}$  in time  $\Delta t$  under influence of force  $\mathbf{F}$ .  
Say its velocity changes from  $\mathbf{v}(t)$  to  $\mathbf{v}(t + \Delta t)$ . Since

$$\mathbf{v}(t + \Delta t) = \mathbf{v}(t) + \frac{d\mathbf{v}}{dt} \Delta t + O(2)$$



# Work

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Then, the change in the kinetic energy is

$$\begin{aligned} & \frac{1}{2}m\mathbf{v}(t + \Delta t) \cdot \mathbf{v}(t + \Delta t) - \frac{1}{2}m\mathbf{v}(t) \cdot \mathbf{v}(t) \\ = & \frac{1}{2}m \frac{d\mathbf{v}}{dt} \cdot \mathbf{v} \Delta t + \frac{1}{2}m\mathbf{v} \cdot \frac{d\mathbf{v}}{dt} \Delta t + O(2) \\ = & \mathbf{F} \cdot \Delta \mathbf{r} \end{aligned}$$

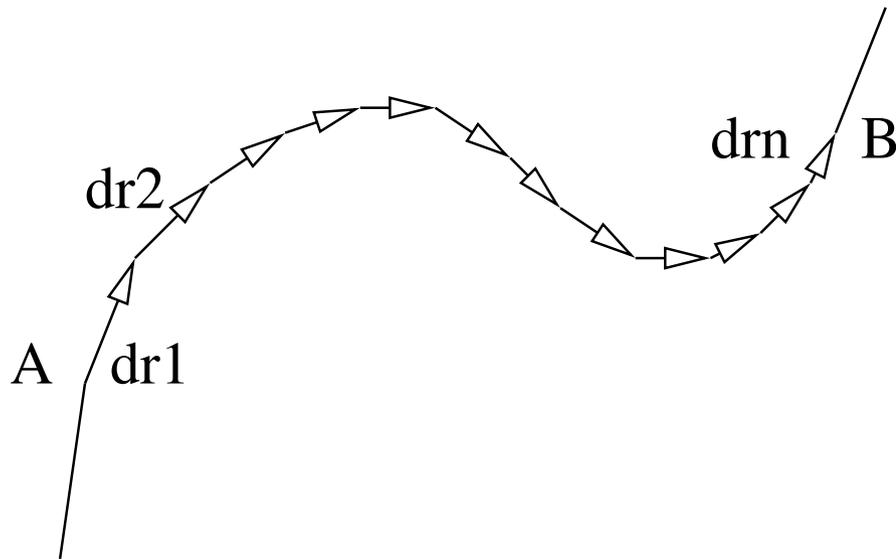
The quantity  $\mathbf{F} \cdot \Delta \mathbf{r}$  is called the work done by force  $\mathbf{F}$  on the particle.

# Work

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Work done by the force  $F$  in moving the particle from position  $A$  to position  $B$  along a given path is given by

$$\begin{aligned}W_{AB} &= \sum \mathbf{F}_i \cdot \Delta \mathbf{r}_i \\ &= \int_A^B \mathbf{F} \cdot d\mathbf{r}\end{aligned}$$



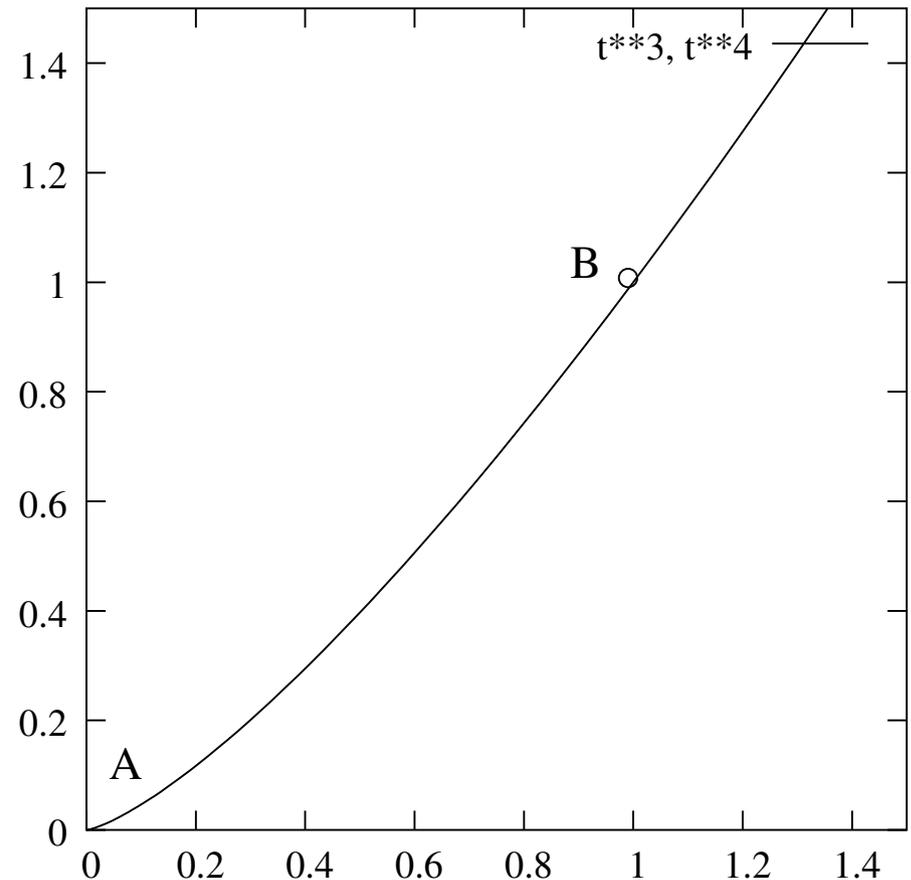
# Example

A particle (1 kg) starts from origin with zero velocity under the influence of force  $\mathbf{F} = 6t\hat{\mathbf{i}} + 12t^2\hat{\mathbf{j}}$ . (SI units.) Then acceleration, velocity and position are

$$\mathbf{a} = 6t\hat{\mathbf{i}} + 12t^2\hat{\mathbf{j}}$$

$$\mathbf{v} = 3t^2\hat{\mathbf{i}} + 4t^3\hat{\mathbf{j}}$$

$$\mathbf{r} = t^3\hat{\mathbf{i}} + t^4\hat{\mathbf{j}}$$



# Example Continued...

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At time  $t$ , In time  $dt$ , suppose that the particle moves by  $d\mathbf{r} = dx\hat{\mathbf{i}} + dy\hat{\mathbf{j}}$ .  
Since  $x = t^3$ , clearly  $dx = 3t^2 dt$  and similarly  $dy = 4t^3 dt$ .

$$d\mathbf{r} = \left(3t^2\hat{\mathbf{i}} + 4t^3\hat{\mathbf{j}}\right) dt$$

Then the work done by the force

$$\begin{aligned}W_{AB} &= \int_A^B \mathbf{F} \cdot d\mathbf{r} \\&= \int_0^1 (18t^3 + 48t^5) dt \\&= \frac{25}{2}\end{aligned}$$

# Properties of Work

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- If there are two forces, say  $\mathbf{F}_1$  and  $\mathbf{F}_2$  acting on a particle, net work done on a particle is, sum of work done by individual forces.

$$\begin{aligned}W_{AB} &= \int_A^B (\mathbf{F}_1 + \mathbf{F}_2) \cdot d\mathbf{r} \\ &= \int_A^B \mathbf{F}_1 \cdot d\mathbf{r} + \int_A^B \mathbf{F}_2 \cdot d\mathbf{r} \\ &= W_{AB}^1 + W_{AB}^2\end{aligned}$$

# Properties of Work

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- If a particle is moved from A to B and then B to C. Then

$$\begin{aligned}W_{AC} &= \int_A^C \mathbf{F} \cdot d\mathbf{r} \\ &= \int_A^B \mathbf{F} \cdot d\mathbf{r} + \int_B^C \mathbf{F} \cdot d\mathbf{r} \\ &= W_{AB} + W_{BC}\end{aligned}$$

- $W_{AB} = -W_{BA}$

# Work-Energy Theorem

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In each time period  $\Delta t$ , the particle moves a distance  $\Delta \mathbf{r}$ . Then the change in Kinetic Energy  $\Delta K$  is related to

$$\Delta K = \mathbf{F} \cdot \Delta \mathbf{r}$$

When particle moves from position A to B, all small contributions add to give

$$\begin{aligned} \frac{1}{2}mv_B^2 - \frac{1}{2}mv_A^2 &= \int_A^B \mathbf{F} \cdot d\mathbf{r} \\ &= W_{AB} \end{aligned}$$

# Work-Energy Theorem

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In case, there are several forces acting on the particle,

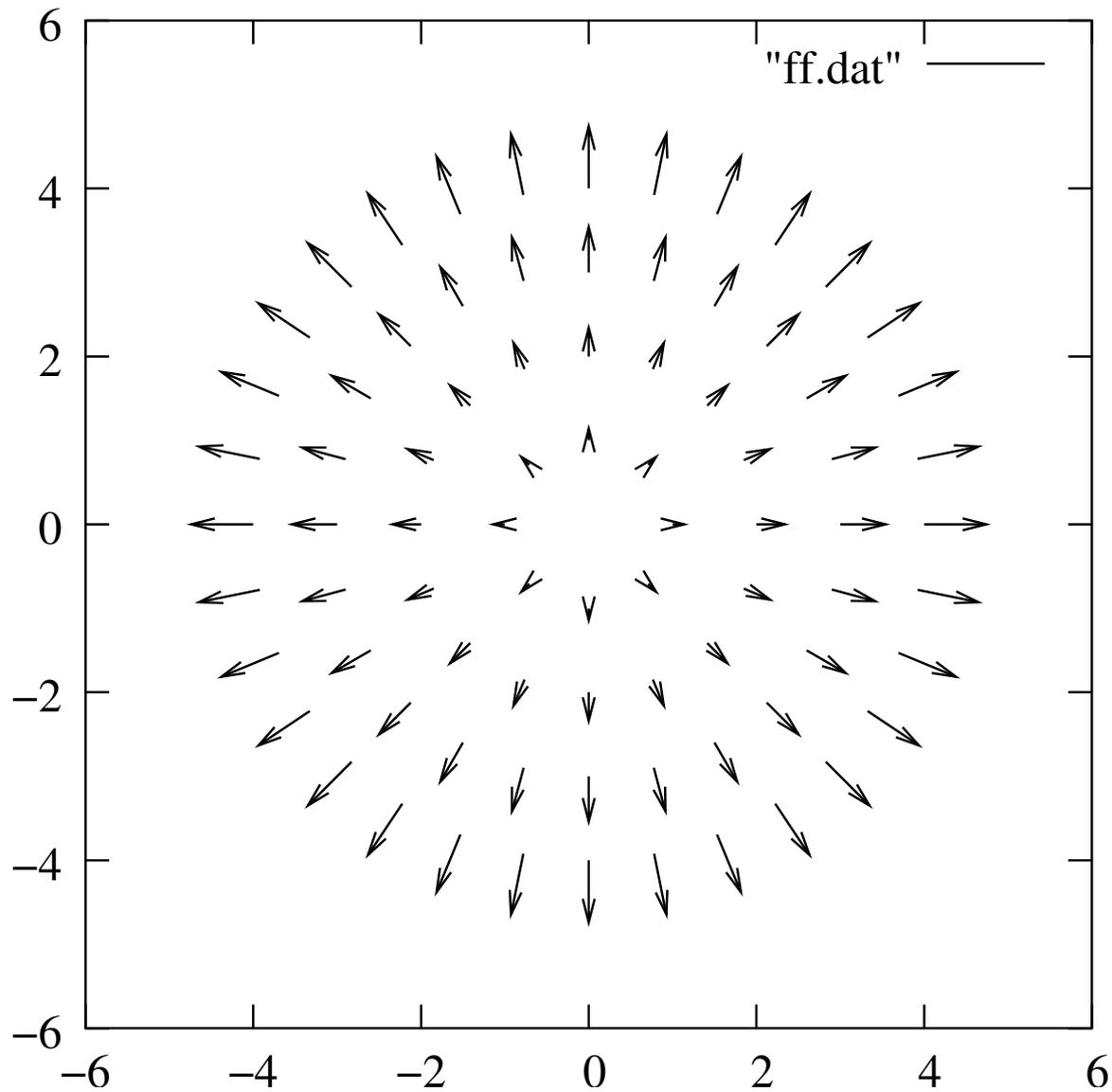
$$\begin{aligned}\frac{1}{2}mv_B^2 - \frac{1}{2}mv_A^2 &= \sum_i \int_A^B \mathbf{F}_i \cdot d\mathbf{r} \\ &= \sum_i W_{AB}^i\end{aligned}$$

# Force Fields

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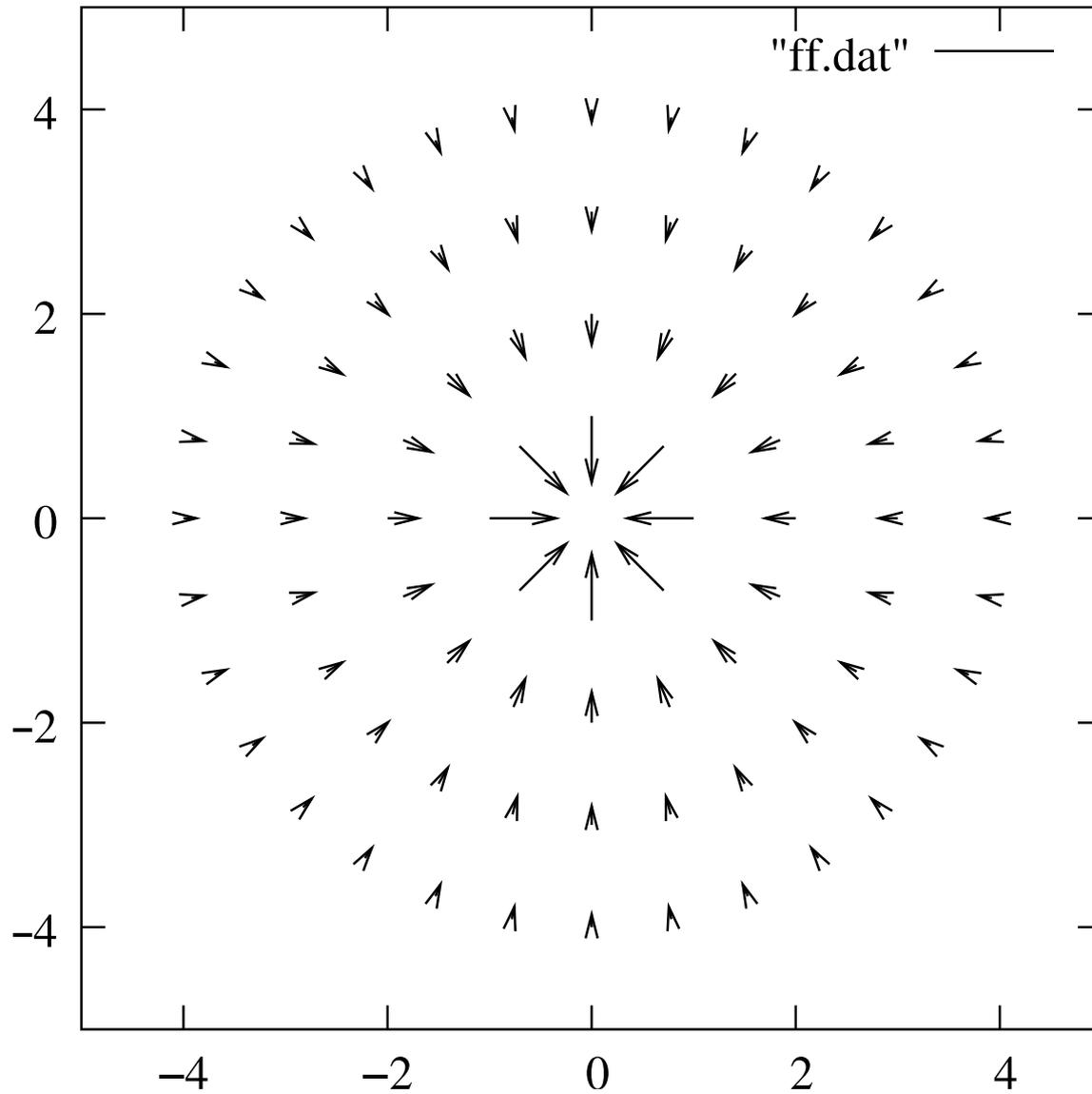
- If a force is a function of position vector only and is defined at all points of the space, then such forces will be called force fields.
- The force fields do not explicitly depend on time.
- The value of the force on a particle at a point does not depend on the history of the particle.

# Force Fields: Examples



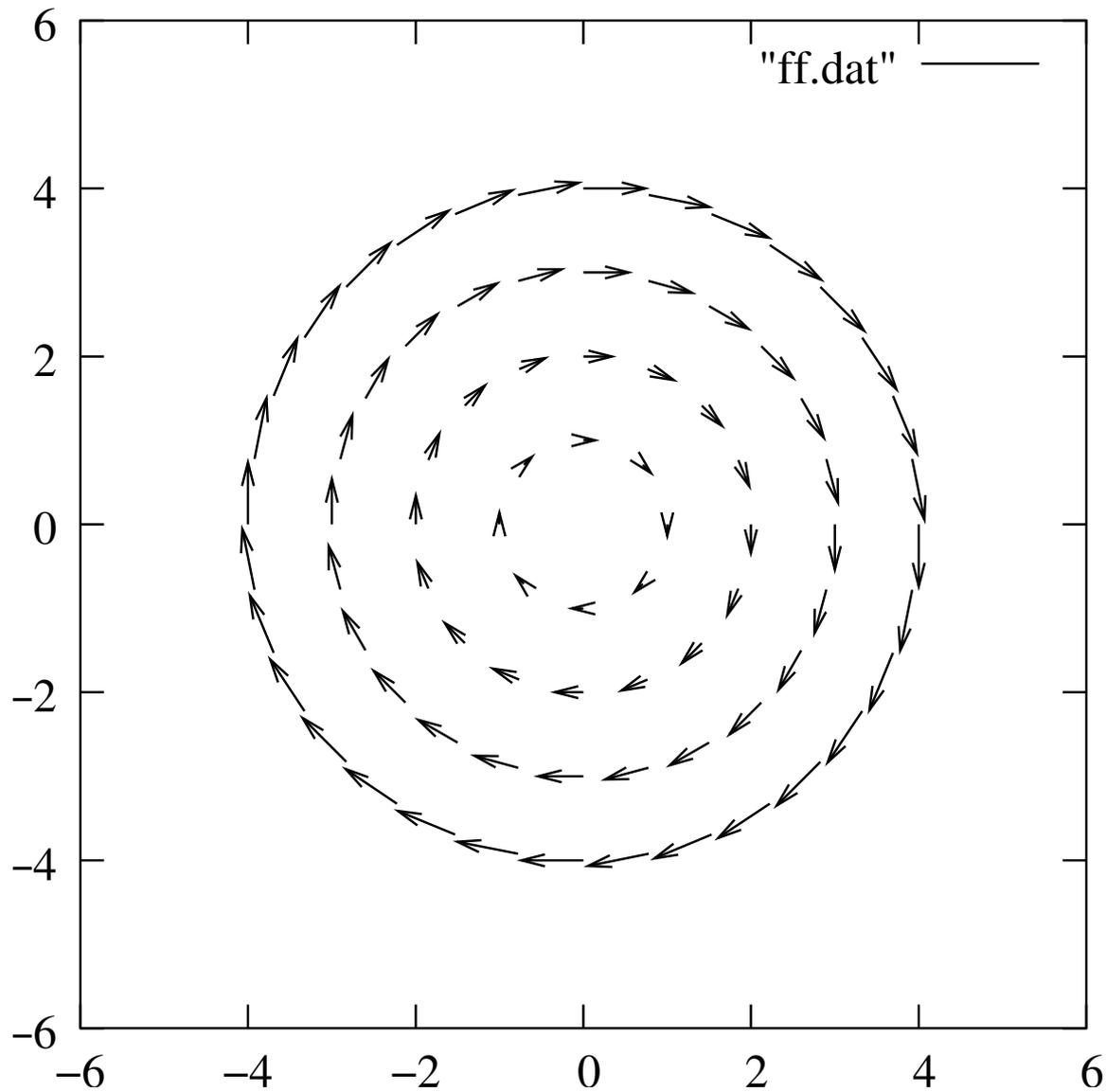
$$\mathbf{F} = x\hat{i} + y\hat{j}$$

# Force Fields: Examples



$$\mathbf{F} = \frac{1}{r^2} \hat{\mathbf{r}}$$

# Force Fields: Examples

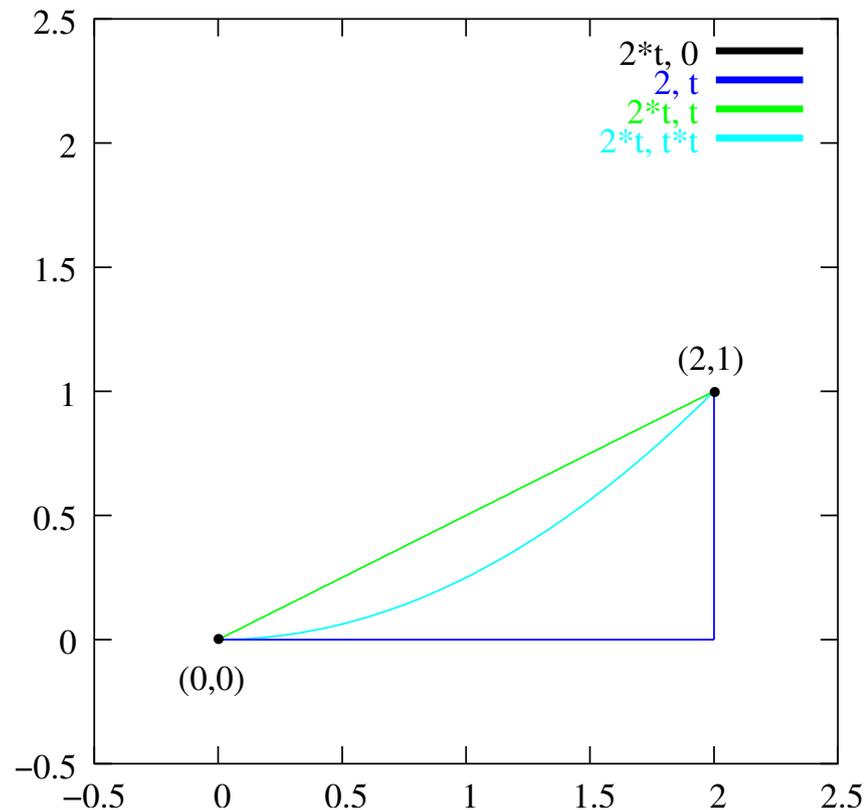


$$\mathbf{F} = y\hat{\mathbf{i}} - x\hat{\mathbf{j}}$$

# Work Done by Force Fields

Let  $\mathbf{F} = x\hat{\mathbf{i}} + y\hat{\mathbf{j}}$ .

A particle moves from origin to  $(2,1)$ . Some possible paths are shown in the figure.



# Examples Continued...

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Given  $\mathbf{F} = x\hat{\mathbf{i}} + y\hat{\mathbf{j}}$ .

Path 1:  $(0, 0) \longrightarrow (2, 0) \longrightarrow (2, 1)$

Along  $(0, 0) \longrightarrow (2, 0)$ ,  $d\mathbf{r} = dx\hat{\mathbf{i}}$

$$\int \mathbf{F} \cdot d\mathbf{r} = \int_0^2 F_x dx = 2$$

Along  $(2, 0) \longrightarrow (2, 1)$ ,  $d\mathbf{r} = dy\hat{\mathbf{j}}$

$$\int \mathbf{F} \cdot d\mathbf{r} = \int_0^1 F_y dy = 1/2$$

Hence work done by  $\mathbf{F}$  along path 1 is

$$W_{AB} = 2 + 1/2 = 2.5$$

# Examples Continued...

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Given  $\mathbf{F} = x\hat{\mathbf{i}} + y\hat{\mathbf{j}}$ .

Path 2:  $y = \frac{1}{2}x$

Along path 2,  $dy = 1/2dx$  and  $d\mathbf{r} = dx\hat{\mathbf{i}} + dy\hat{\mathbf{j}} = (\hat{\mathbf{i}} + 1/2\hat{\mathbf{j}})dx$

$$\begin{aligned}\int \mathbf{F} \cdot d\mathbf{r} &= \int_0^2 \mathbf{F} \cdot (\hat{\mathbf{i}} + 1/2\hat{\mathbf{j}})dx \\ &= \int_0^2 (x + (x/2)(1/2))dx \\ &= 2.5\end{aligned}$$

Hence work done by  $\mathbf{F}$  along path 2 is

$$W_{AB} = 2.5$$

# Examples Continued...

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Given  $\mathbf{F} = x\hat{\mathbf{i}} + y\hat{\mathbf{j}}$ .

Path 3:  $y = \frac{1}{4}x^2$

Along path 3,  $dy = \frac{1}{2}x dx$  and  $d\mathbf{r} = dx\hat{\mathbf{i}} + dy\hat{\mathbf{j}} = (\hat{\mathbf{i}} + (x/2)\hat{\mathbf{j}})dx$

$$\begin{aligned}\int \mathbf{F} \cdot d\mathbf{r} &= \int_0^2 \mathbf{F} \cdot (\hat{\mathbf{i}} + (x/2)\hat{\mathbf{j}})dx \\ &= \int_0^2 (x + (x^2/4)(x/2))dx \\ &= 2.5\end{aligned}$$

Hence work done by  $\mathbf{F}$  along path 3 is

$$W_{AB} = 2.5$$

# Another Example

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Given  $\mathbf{F} = y\hat{\mathbf{i}} - x\hat{\mathbf{j}}$ .

Path 1:  $(0, 0) \longrightarrow (1, 0) \longrightarrow (1, 1)$

Along  $(0, 0) \longrightarrow (1, 0)$ ,  $d\mathbf{r} = dx\hat{\mathbf{i}}$  and  $\mathbf{F} = -x\hat{\mathbf{j}}$

$$\int \mathbf{F} \cdot d\mathbf{r} = \int_0^1 F_x dx = 0$$

Along  $(1, 0) \longrightarrow (1, 1)$ ,  $d\mathbf{r} = dy\hat{\mathbf{j}}$  and  $\mathbf{F} = y\hat{\mathbf{i}} - \hat{\mathbf{j}}$

$$\int \mathbf{F} \cdot d\mathbf{r} = \int_0^1 F_y dy = -1$$

Hence work done by  $\mathbf{F}$  along path 1 is

$$W_{AB} = 0 - 1 = -1$$

# Another Example

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Given  $\mathbf{F} = y\hat{\mathbf{i}} - x\hat{\mathbf{j}}$ .

Path 2:  $(0, 0) \longrightarrow (0, 1) \longrightarrow (1, 1)$

Along  $(0, 0) \longrightarrow (0, 1)$ ,  $d\mathbf{r} = dy\hat{\mathbf{j}}$  and  $\mathbf{F} = y\hat{\mathbf{i}}$

$$\int \mathbf{F} \cdot d\mathbf{r} = \int_0^1 F_y dy = 0$$

Along  $(0, 1) \longrightarrow (1, 1)$ ,  $d\mathbf{r} = dx\hat{\mathbf{i}}$  and  $\mathbf{F} = \hat{\mathbf{i}} - x\hat{\mathbf{j}}$

$$\int \mathbf{F} \cdot d\mathbf{r} = \int_0^1 F_x dx = 1$$

Hence work done by  $\mathbf{F}$  along path 1 is

$$W_{AB} = 0 + 1 = 1$$

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Work done along path 1 and path 2 are different!